#### Black Hole Spins following Minor Mergers

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Distribution of black hole spins following a sequence of minor mergers

- A compact object (CO) is captured by a more massive black hole (BH)
- Radiation reaction leads to CO inspiral
- CO angular momentum and mass are added to the black hole's
- If a BH grows exclusively through a sequence of minor mergers, what will be its spin?

#### Spin evolution via minor mergers

- Only orbital angular momentum of compact object at LSO is relevant, not its spin
- Inclination is ~constant over inspiral and isotropically distributed

$$\cos \iota = \frac{L_z}{\sqrt{L_z^2 + Q}}$$

• Total angular momentum for small  $\chi$  is

$$\sqrt{L_z^2+Q}\approx Mm\sqrt{12}\left[1-\frac{1}{2}\left(\frac{2}{3}\right)^{3/2}\chi\cos\iota\right]$$

• Minor merger of mass *m* object with mass *M* Kerr black hole changes spin parameter  $\chi = S/M^2$  to

$$\chi' \approx \frac{1}{(M+m)^2} \sqrt{(\chi M^2 + L_z)^2 + Q}$$

## Monte Carlo simulations of spin distribution

• Evolution from M=5m to M=10m, starting with  $\chi$ =0.1 and  $\chi$ =0.9 (see also [Miller, 2002])



### Fokker-Planck equation for the evolution of the spin distribution

$$\frac{\partial}{\partial t}f(x,t) = -\frac{\partial}{\partial x}\left[\mu(x,t)f(x,t)\right] + \frac{1}{2}\frac{\partial^2}{\partial x^2}\left[\sigma^2(x,t)f(x,t)\right]$$

- [Hughes & Blandford, 2004] derived a 3-d Fokker-Planck equation but did not solve it
- We parametrize mass by a "time" parameter t=M/m; when  $\chi$ t>>1,

$$\mu(\chi,t) = \frac{\langle d\chi \rangle}{dt} = \frac{\chi}{t} \left( -2 - \frac{4\sqrt{2}}{9} \right) + \frac{4}{\chi t^2}$$

$$\sigma^{2}(\chi, t) = \frac{\langle (d\chi)^{2} \rangle}{dt} = \frac{4}{t^{2}} \left(1 + \frac{4\sqrt{2}\chi^{2}}{9} - \chi^{2}\right)$$

• The Fokker-Planck equation still looks complicated...

$$\frac{\partial}{\partial t}f(\chi,t) = -\frac{\partial}{\partial \chi} \left[ \frac{\chi}{t} \left( -2 - \frac{4\sqrt{2}}{9} + \frac{4}{\chi^2 t} \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi,t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[ \frac{4}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) \right] + \frac{1}{2} \frac{2}{t^2} \left( 1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) \right]$$

# Analytical approximation to evolution of probability distribution for the black hole spin

- Probability distribution is roughly Gaussian, so consider evolution of mean and standard deviation
- If  $\chi^2 t >> 1$ , then

$$\approx \bar{\chi}_0 \left(\frac{t}{t_0}\right)^a \approx \bar{\chi}_0 \left(\frac{M_0}{M}\right)^{2.63}$$
 where

$$a \equiv -2 - 4\sqrt{2}/9 \approx -2.63$$

- The random walk is damped, since retrograde orbits have a greater angular momentum at LSO, otherwise exponent would be -2; [Hughes & Blandford found an exponent of -2.4]
- As long as  $\chi t >> 1$ , mean spin tends toward

 $\bar{\chi}$ 

$$\bar{\chi} \rightarrow \sqrt{\frac{4}{-at}} \approx \sqrt{\frac{1.5}{t}}$$

• As long as  $\chi t >> 1$ , spin standard deviation tends toward



where 
$$b \equiv 4\sqrt{2}/9 - 1$$

## Spin distribution for intermediate mass black holes

- Evolution from t=M/m=50 to t=100 (e.g., from M=70 to M=140 solar masses via capture of m=1.4 solar-mass NSs)
- If initial  $\chi$ =0.1, then mean spin at t=100 is 0.162,  $\sigma$ =0.066
- If initial  $\chi$ =0.9, then mean spin at t=100 is 0.233,  $\sigma$ =0.087



## Effect of spin on detection range for LIGO IMRIs

• Prograde inspirals can be seen further (higher LSO frequency)



#### Advanced LIGO:

$$rac{\mathrm{Rate}_{\mathrm{spin}}}{\mathrm{Rate}_{\mathrm{no-spin}}} \sim 1 + 2\chi^2 \left(rac{M}{100~M_{\odot}}
ight)$$

## Effect of spin on detection range for LISA EMRIs

- LISA mission duration limits visible portion of inspiral
- Redshift must be taken into account



M=10<sup>6</sup> M<sub>sun</sub>, m=10 M<sub>sun</sub>, SNR=30

 $M=10^7 M_{sun}$ , m=10  $M_{sun}$ , SNR=30

#### Summary

- Found expressions for BH spin distribution following minor mergers
- If intermediate-mass black holes in globular clusters grow via minor mergers, they will have low spins
- Central black hole spins can aid in detection of gravitational waves from intermediate- or extrememass-ratio inspirals
- This effect can cause a bias in favor of detecting inspirals into rapidly spinning BHs